

Testing Higgs Self-Couplings at High-Energy Linear Colliders

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Abstract. In order to verify the Higgs mechanism experimentally, the Higgs self-couplings have to be probed. These couplings allow the reconstruction of the characteristic Higgs potential responsible for the electroweak symmetry breaking. The couplings are accessible in a variety of multiple Higgs production processes. The theoretical analysis including the most relevant channels for the production of neutral Higgs boson pairs at high-energy and high-luminosity e^+e^- linear colliders will be presented in this note.

1. Within the Higgs mechanism [1] the electroweak gauge bosons and fundamental matter particles acquire their masses through the interaction with a scalar field. The self-interaction of the scalar field induces, via a non-vanishing field strength $v = (\sqrt{2}G_F)^{-1/2} \approx 246$ GeV, the spontaneous breaking of the electroweak $SU(2)_L \times U(1)_Y$ symmetry down to the electromagnetic $U(1)_{EM}$ symmetry.

To establish the Higgs mechanism experimentally, the self-energy potential of the Standard Model [2], $V = \lambda (\Phi^\dagger \Phi - v^2/2)^2$, with a minimum at $\langle \Phi \rangle_0 = v/\sqrt{2}$ must be reconstructed. This task requires the measurement of the Higgs self-couplings of the physical Higgs boson H , which can be read off directly from the potential

$$V = \frac{M_H^2}{2} H^2 + \frac{M_H^2}{2v} H^3 + \frac{M_H^2}{8v^2} H^4 \quad (1)$$

As evident from Eq. (1), in the SM the trilinear and quadrilinear vertices are uniquely determined by the mass of the Higgs boson, $M_H = \sqrt{2\lambda}v$.

The trilinear self-coupling $\lambda = 6\sqrt{2}\lambda$ in units of $v/\sqrt{2}$ is accessible directly in Higgs pair production at high-energy e^+e^- linear colliders. For c.m. energies up to about 1 TeV, double Higgs-strahlung [3,4]

$$e^+e^- \rightarrow ZHH \quad (2)$$

is the most promising process [5]. The process includes the amplitude involving the trilinear Higgs self-coupling and two additional amplitudes due to the standard electroweak gauge interactions, cf. Fig. 1, so that it is a binomial in λ_{HHH} .

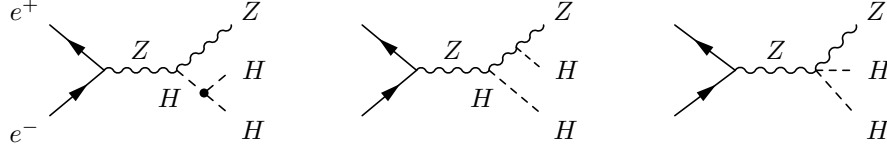


FIGURE 1. Subprocesses contributing to double Higgs-strahlung, $e^+e^- \rightarrow ZHH$, in the Standard Model at e^+e^- linear colliders.

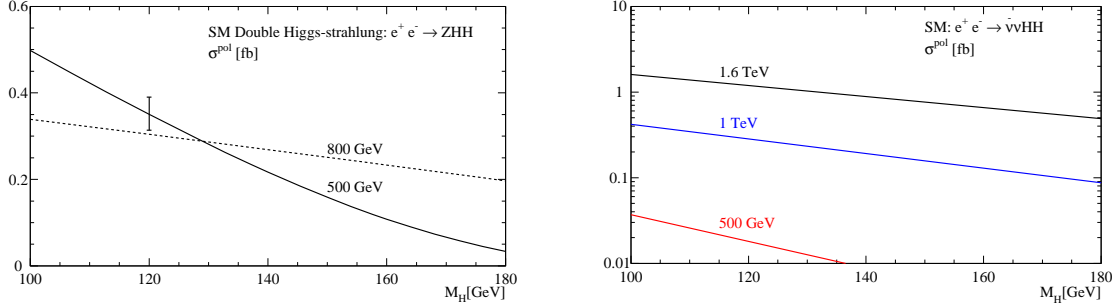


FIGURE 2. (a) The cross section for double Higgs-strahlung, $e^+e^- \rightarrow ZHH$, in the Standard Model at two collider energies: $\sqrt{s} = 500$ GeV and 800 GeV. The electron/positron beams are taken oppositely polarized. The vertical bar corresponds to a variation of the trilinear Higgs coupling between 0.8 and 1.2 of the SM value. (b) The cross section for WW double Higgs fusion, $e^+e^- \rightarrow \nu_e \bar{\nu}_e HH$, at $\sqrt{s} = 500$ GeV, 1 TeV and 1.6 TeV. The initial beams are polarized.

As evident from Fig. 2a the cross section is very sensitive to the trilinear Higgs self-coupling and amounts up to 0.35 fb for $M_H = 120$ GeV and a c.m. energy of 500 GeV. Scaling with the energy, it decreases to 0.3 fb at $\sqrt{s} = 800$ GeV. Experimental detector simulations of signal and background processes in the SM have demonstrated that the Higgs self-coupling can be extracted with an accuracy of $\sim 20\%$ for $M_H = 120$ GeV at high luminosity $\int \mathcal{L} = 2 \text{ ab}^{-1}$ [6].

The WW double Higgs fusion process [4,7]

$$e^+e^- \rightarrow \nu_e \bar{\nu}_e HH \quad (3)$$

which increases with rising \sqrt{s} , can be exploited for larger energies, cf. Fig. 2b. [$\sigma = 0.37$ fb for $M_H = 120$ GeV and $\sqrt{s} = 1$ TeV, polarized e^+e^- beams.]

Triple Higgs production is sensitive to the quadrilinear Higgs self-coupling. Due to the suppressed coupling and an additional particle in the final state, the cross section $\sigma(e^+e^- \rightarrow ZHHH)$ is only of $\mathcal{O}(\text{ab})$ and therefore not measurable at typical linear collider energies and luminosities [5].

2. In the Minimal Supersymmetric Extension of the Standard Model (MSSM) with five physical Higgs particles h, H, A and H^\pm [8], a plethora of trilinear and quadrilinear Higgs self-couplings can be realized. The CP-invariant couplings

λ	double Higgs–strahlung				triple Higgs–production			
	Zhh	ZHh	ZHH	ZAA	Ahh	AHh	AHH	AAA
hhh	×				×			
Hhh	×	×			×	×		
HHh		×	×			×	×	
HHH			×				×	
hAA				×	×	×		×
HAA				×		×	×	×

TABLE 1: The trilinear couplings between neutral CP-even and CP-odd MSSM Higgs bosons, which can generically be probed in double Higgs-strahlung and associated triple Higgs-production, are marked by a cross. [The matrix for WW fusion is isomorphic to the matrix for Higgs-strahlung.]

among the neutral Higgs bosons, $\lambda_{hhh}, \lambda_{Hhh}, \lambda_{HHh}, \lambda_{HHH}, \lambda_{hAA}, \lambda_{HAA}$, are involved in a large number of processes [5,9]. The double and triple Higgs production processes and the trilinear couplings, that can be probed in the respective process, are listed in Table 1. The system is solvable for all λ 's up to discrete ambiguities. However, in practice, not all processes are accessible experimentally so that one has to follow the reverse direction in this case: Comparing the theoretical predictions with the experimental results of the accessible channels, the trilinear Higgs self-couplings can be tested stringently.

The process $e^+e^- \rightarrow Zhh$ is sensitive to the trilinear coupling of the light CP-even Higgs boson h ,

$$\lambda_{hhh} = 3 \cos 2\alpha \sin(\beta + \alpha) + \mathcal{O}(G_F M_t^4 / M_Z^2) \quad (4)$$

expressed in the mixing angles α and β , in a large range of the MSSM parameter space, as can be inferred from Fig. 3. It shows the 2σ sensitivity area in the $[M_A, \tan \beta]$ plane for a non-zero coupling at an integrated luminosity of 2 ab^{-1} . The cross section is required to exceed 0.01 fb . The sensitivity areas are significantly smaller for processes involving heavy Higgs bosons H and A in the final state. Details can be found in Ref. [5].

3. In summary. The measurement of the Higgs self-couplings is essential for the reconstruction of the characteristic self-energy potential. The large luminosities, which are available at future high-energy e^+e^- linear colliders, allow the measurement of the trilinear Higgs self-couplings via double Higgs-strahlung and WW double Higgs fusion.

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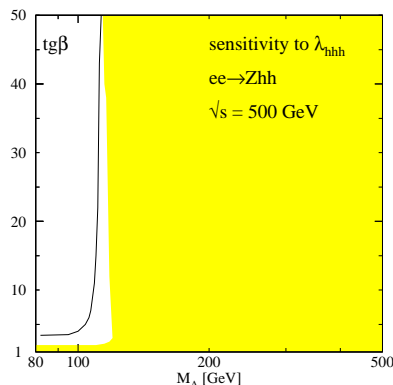


FIGURE 3. MSSM: Sensitivity to the coupling λ_{hhh} of the light CP-even neutral Higgs boson h in the process $e^+e^- \rightarrow Zhh$ for a collider energy of 500 GeV (no mixing).

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REFERENCES

1. P.W. Higgs, Phys. Lett. **12** (1964) 132; and Phys. Rev. **145** (1966) 1156; F. Englert and R. Brout, Phys. Rev. Lett. **13** (1964) 321; G.S. Guralnik, C.R. Hagen and T.W. Kibble, Phys. Rev. Lett. **13** (1964) 585.
2. For a review see: J.F. Gunion, H.E. Haber, G. Kane and S. Dawson, *The Higgs Hunter's Guide*, Addison-Wesley, 1990; A. Djouadi, Int. J. Mod. Phys. **A10** (1995) 1; M. Spira and P.M. Zerwas, *Int. Universitätswochen*, Schladming 1997, hep-ph/9803257; C. Quigg, Acta Phys. Polon. **B30** (1999) 2145-2192.
3. G. Gounaris, D. Schildknecht and F. Renard, Phys. Lett. **B83** (1979) 191 and (E) **B89** (1980) 437; V. Barger, T. Han and R.J.N. Phillips, Phys. Rev. **D38** (1988) 2766.
4. V.A. Ilyin, A.E. Pukhov, Y. Kurihara, Y. Shimizu and T. Kaneko, Phys. Rev. **D54** (1996) 6717.
5. A. Djouadi, W. Kilian, M. Mühlleitner and P.M. Zerwas, Eur. Phys. J. **C10** (1999) 27 and *ibid.* DESY 99-171, PM/99-55, TTP99-48, hep-ph/0001169; M.M. Mühlleitner, PhD thesis, DESY-THESIS-2000-033, hep-ph/0008127.
6. C. Castanier, P. Gay, P. Lutz and J. Orloff, LC-PHSM-2000-061, hep-ex/0101028.
7. F. Boudjema and E. Chopin, Z. Phys. **C73** (1996) 85; V. Barger and T. Han, Mod. Phys. Lett. **A5** (1990) 667; D.A. Dicus, K.J. Kallianpur and S.S.D. Willenbrock, Phys. Lett. **B200** (1988) 187; A. Abbasabadi, W.W. Repko, D.A. Dicus and R. Vega, Phys. Rev. **D38** (1988) 2770 and *ibid.* Phys. Lett. **B213** (1988) 386; A. Dobrovolskaya and V. Novikov, Z. Phys. **C52** (1991) 427.
8. J.F. Gunion and H.E. Haber, Nucl. Phys. **B272** (1986) 1 and **B278** (1986) 449.
9. A. Djouadi, H.E. Haber and P.M. Zerwas, Phys. Lett. **B375** (1996) 203 and (E); P. Osland and P.N. Pandita, Phys. Rev. **D59** (1999) 055013.